# Comments on non-isometric T-duality 

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## Outline

(1) Review of Isometric T-duality
(2) Non-isometric T-duality
(3) Equivalence
(4) Examples

## Setting up notation

Consider a non-linear sigma model $X: \Sigma \rightarrow M$ described by the following action:

$$
S=\int_{\Sigma} g_{i j} d X^{i} \wedge \star d X^{j}+\int_{\Sigma} B_{i j} d X^{i} \wedge d X^{j}
$$

In this talk we will ignore the dilaton, and assume that both $g$ and $B$ are globally defined fields on $M$.

## Gauging isometries

Suppose now that there are vector fields generating the following global symmetry:

$$
\delta_{\epsilon} X^{i}=v_{a}^{i} \epsilon^{a}
$$

for $\epsilon^{a}$ constant. The sigma model action is invariant under this transformation if

$$
\mathcal{L}_{V_{a}} g=0 \quad \mathcal{L}_{V_{a}} B=0
$$

If this is the case, we can gauge the model by promoting the global symmetry to a local one.

## The gauged action

Introducing gauge fields $A^{a}$ and Lagrange multipliers $\eta_{a}$, the gauged action is

$$
S_{G}=\int_{\Sigma} g_{i j} D X^{i} \wedge \star D X^{j}+\int_{\Sigma} B_{i j} D X^{i} \wedge D X^{j}+\int_{\Sigma} \eta_{a} F^{a}
$$

where

- $F=d A+A \wedge A$ is the standard Yang-Mills field strength
- $D X^{i}=d X^{i}-v_{a}^{i} A^{a}$ are the gauge covariant derivatives.


## Gauge invariance

The gauged action is invariant with respect to the following (local) gauge transformations:

$$
\begin{aligned}
\delta_{\epsilon} X^{i} & =v_{a}^{i} \epsilon^{a} \\
\delta_{\epsilon} A^{a} & =d \epsilon^{a}+C_{b c}^{a} A^{b} \epsilon^{c} \\
\delta_{\epsilon} \eta_{a} & =-C_{a b}^{c} \epsilon^{b} \eta_{c}
\end{aligned}
$$

## T-duality



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On the other hand, we can eliminate the non-dynamical gauge fields $A$, obtaining the dual sigma model.

## Can we do it without isometries?

The existence of global symmetries is a very stringent requirement. A generic metric will not have any Killing vectors.

## Question

Is it possible to follow the same procedure when the vector fields are not Killing vectors?

## Gauging without isometry

Kotov and Strobl ${ }^{1}$ generalised the standard gauging using Lie algebroids (see Kyle's talk).
Chatzistavrakidis, Deser, and Jonke ${ }^{2}$ applied this non-isometric gauging to the Buscher procedure we just reviewed.

They introduce a matrix-valued one-form $\omega_{a}^{b}$ satisfying

$$
\begin{aligned}
\mathcal{L}_{v_{a}} g & =\omega_{a}^{b} \vee \iota_{v_{b}} g \\
\mathcal{L}_{v_{a}} B & =\omega_{a}^{b} \wedge \iota_{v_{b}} B
\end{aligned}
$$

[^0]
## The gauged action

The gauged action is almost the same:

$$
S_{G}^{\omega}=\int_{\Sigma} g_{i j} D X^{i} \wedge \star D X^{j}+\int_{\Sigma} B_{i j} D X^{i} \wedge D X^{j}+\int_{\Sigma} \eta_{a} F_{\omega}^{a}
$$

where the curvature is now given by

$$
F_{\omega}^{a}=d A^{a}+\frac{1}{2} C_{b c}^{a}(X) A^{b} \wedge A^{c}-\omega_{b i}^{a} A^{b} \wedge D X^{i}
$$

## Modified gauge invariance

The modified gauge transformations are now

$$
\begin{aligned}
\delta_{\epsilon} X^{i} & =v_{a}^{i} \epsilon^{a} \\
\delta_{\epsilon} A^{a} & =d \epsilon^{a}+C_{b c}^{a} A^{b} \epsilon^{c}+\omega_{b i}^{a} \epsilon^{b} D X^{i} \\
\delta_{\epsilon} \eta_{a} & =-C_{a b}^{c} \epsilon^{b} \eta_{c}+v_{a}^{i} \omega_{b i}^{c} \epsilon^{b} \eta_{c}
\end{aligned}
$$

## T-duality



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In principle, we could use this to construct T-duals of spaces which have no isometries.

## The problem?

This proposal is equivalent to non-abelian T-duality. ${ }^{3}$
That is, if we can find a set of vector fields and $\omega_{a}^{b}$ which give a non-isometric T-dual, then there exists a set of Killing vectors for the model. The T-dual with respect to these Killing vectors is the same as the non-isometric T-dual.

## A necessary condition for gauge invariance

Gauge invariance of the action requires the structure functions to be constant, as well as the vanishing of the following variation:

$$
\delta_{\epsilon}\left(\eta_{a} F_{\omega}^{a}\right)=\eta_{a}\left(d \omega_{b}^{a}+\omega_{c}^{a} \wedge \omega_{b}^{c}\right) \epsilon^{b}+\mathcal{O}(A)+\mathcal{O}\left(A^{2}\right) .
$$

We therefore require that $\omega_{a}^{b}$ is flat:

$$
R_{a}^{b}=d \omega_{a}^{b}+\omega_{c}^{b} \wedge \omega_{a}^{c}=0
$$

and this tells us that $\omega_{a}^{b}$ is of the form $K^{-1} d K$ for some $K_{a}^{b}(X)$.

## A field redefinition

Using this $K$, we can perform the following field redefinitions:

$$
\begin{aligned}
\widehat{A}^{a} & =K_{b}^{a} A^{b} \\
\widehat{\eta}^{a} & =\eta_{b}\left(K^{-1}\right)_{a}^{b} \\
\widehat{v}_{a} & =v_{b}^{i}\left(K^{-1}\right)_{a}^{b}
\end{aligned}
$$

## The non-abelian action!

The gauged action can now be rewritten in terms of the new fields $\left(X^{i}, \widehat{A}^{a}, \widehat{\eta}_{a}\right)$.

$$
\begin{aligned}
S_{G}^{\omega}[X, \widehat{A}, \widehat{\eta}] & =\int_{\Sigma} g_{i j} \widehat{D X}^{i} \wedge \star \widehat{D X}^{j}+\int_{\Sigma} B_{i j} \widehat{D X}^{i} \wedge \widehat{D X}^{j}+\int_{\Sigma} \widehat{\eta}_{a} \widehat{F}^{a} \\
& =S_{G}[X, \widehat{A}, \widehat{\eta}]
\end{aligned}
$$

where

$$
\widehat{F}^{a}=d \widehat{A}^{a}+\frac{1}{2} \widehat{C}_{b c}^{a} \widehat{A}^{b} \wedge \widehat{A}^{c}
$$

The gauge transformations become the usual non-abelian gauge transformations, and a short computation reveals

$$
\mathcal{L}_{\widehat{V}_{a}} g=0 \quad \mathcal{L}_{\widehat{V}_{a}} B=0
$$

## Conclusion

This proposal is equivalent, via a field redefinition, to the standard non-abelian T-duality

## Examples!

## First example

Consider the 3D Heisenberg Nilmanifold, or twisted torus. It has a metric given by

$$
d s^{2}=d x^{2}+(d y-x d z)^{2}+d z^{2}
$$



The non-abelian T-dual of this space is given by

$$
\begin{aligned}
\widehat{d s^{2}} & =d Y^{2}+\frac{1}{1+Y^{2}}\left(d X^{2}+d Z^{2}\right) \\
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## First example

We can gain a better understanding of the geometry by writing the manifold as a group:

$$
\text { Heis }:=\left\{\left(\begin{array}{ccc}
1 & x & y \\
0 & 1 & z \\
0 & 0 & 1
\end{array}\right): x, y, z \in \mathbb{R}\right\}
$$

(left-invariant) MC forms $=(d x, d y-x d z, d z)$
(right-invariant) vector fields $=\left(\partial_{x}+z \partial_{y}, \partial_{y}, \partial_{z}\right)$

## First example

We could instead try to gauge this space non-isometrically using the left-invariant vector fields: $\left\{\partial_{x}, \partial_{y}, x \partial_{y}+\partial_{z}\right\}$.

These are not all isometries:

$$
\begin{aligned}
& \mathcal{L}_{V_{1}} g=-d y \otimes d x-d z \otimes d y+2 x d z \otimes d z \\
& \mathcal{L}_{V_{2}} g=0 \\
& \mathcal{L}_{V_{3}} g=d x \otimes d y+d y \otimes d x-x d x \otimes d z-x d z \otimes d x
\end{aligned}
$$

and they don't commute:

$$
\left[v_{1}, v_{3}\right]=v_{2}
$$

however...

## First example ${ }^{4}$

If we take $\omega_{3}^{2}=d x$ and $\omega_{1}^{2}=-d z$, with other components vanishing, the non-isometric gauging constraints are satisfied and we can calculate the non-isometric T-dual model.

$$
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\widehat{d s^{2}} & =d Y^{2}+\frac{1}{1+Y^{2}}\left(d X^{2}+d Z^{2}\right) \\
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[^1]A. Chatzistavrakidis, A. Deser, L. Jonke

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Unsurprisingly, it is also the T-fold.

[^2]
## Second example

Consider $S^{3}$ with the round metric and $B=0$.


This metric has an $S O(4)$ group of isometries, and we can find the non-abelian T-dual with respect to an $S U(2)$ subgroup of this.

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## Second example

We can write the round metric as

$$
g=\lambda^{1} \otimes \lambda^{1}+\lambda^{2} \otimes \lambda^{2}+\lambda^{3} \otimes \lambda^{3}
$$

where the $\lambda^{i}$ are the left-invariant Maurer-Cartan forms.
The right-invariant vector fields are isometries of this metric, so let's try gauging with respect to the left-invariant vector fields ${ }^{5}$.

[^3] non-isometrically

## Second example

The Lie derivatives of the metric with respect to the left-invariant vector fields, $L_{a}$ are

$$
\begin{aligned}
\mathcal{L}_{L_{a}} g & =-\sum_{b} C_{a c}^{b} \lambda^{c} \vee \lambda^{b} \\
& =-C_{a c}^{b} \lambda^{c} \vee \iota_{L_{b}} g
\end{aligned}
$$

We can do non-isometric T-duality by taking $\omega_{a}^{b}=-C_{a c}^{b} \lambda^{c}$.

## Second example

The remaining gauging constraints are satisfied, and we can calculate the non-isometric T-dual. It is the 'cigar' metric, as expected.


## Comments

- The equivalence of non-isometric and non-abelian T-duality remains valid for non-exact $H$
- Geometric interpretation of $\omega_{a}^{b}$ as a connection on a Lie algebroid
- There are proposals for alternate gauging. Unknown how to incorporate into T-duality
- non-flat $\omega_{a}^{b}$
- include a term $\phi_{\text {ai }}^{b} \epsilon^{b} \star D X^{i}$ into $\delta_{\epsilon} A^{a}$


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## Thanks!


[^0]:    ${ }^{1}$ [1403.8119]
    ${ }^{2}$ [1509.01829] and [1604.03739]

[^1]:    ${ }^{4}$ Gauged non-isometrically in [1509:01829]

[^2]:    ${ }^{4}$ Gauged non-isometrically in [1509:01829]
    A. Chatzistavrakidis, A. Deser, L. Jonke

[^3]:    ${ }^{5}$ These also happen to be isometries of the metric, but let's try to gauge them

