### Comments on non-isometric T-duality

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Based on [1705.09254] with **P. Bouwknegt, C. Klimčík**, and **K. Wright** 

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## Outline

Review of Isometric T-duality

#### 2 Non-isometric T-duality





Consider a non-linear sigma model  $X : \Sigma \to M$  described by the following action:

$$S = \int_{\Sigma} g_{ij} dX^i \wedge \star dX^j + \int_{\Sigma} B_{ij} dX^i \wedge dX^j$$

In this talk we will ignore the dilaton, and assume that both g and B are globally defined fields on M.

Suppose now that there are vector fields generating the following global symmetry:

$$\delta_{\epsilon} X^{i} = v^{i}_{a} \epsilon^{a}$$

for  $\epsilon^{\rm a}$  constant. The sigma model action is invariant under this transformation if

$$\mathcal{L}_{v_a}g=0 \qquad \mathcal{L}_{v_a}B=0$$

If this is the case, we can gauge the model by promoting the global symmetry to a local one.

Introducing gauge fields  $A^{\rm a}$  and Lagrange multipliers  $\eta_{\rm a},$  the gauged action is

$$S_{G} = \int_{\Sigma} g_{ij} DX^{i} \wedge \star DX^{j} + \int_{\Sigma} B_{ij} DX^{i} \wedge DX^{j} + \int_{\Sigma} \eta_{a} F^{a}$$

where

- $F = dA + A \land A$  is the standard Yang-Mills field strength
- $DX^i = dX^i v_a^i A^a$  are the gauge covariant derivatives.

The gauged action is invariant with respect to the following (local) gauge transformations:

$$\delta_{\epsilon} X^{i} = v_{a}^{i} \epsilon^{a}$$
  
$$\delta_{\epsilon} A^{a} = d\epsilon^{a} + C_{bc}^{a} A^{b} \epsilon^{c}$$
  
$$\delta_{\epsilon} \eta_{a} = -C_{ab}^{c} \epsilon^{b} \eta_{c}$$





Varying the Lagrange multipliers forces the field strength F to vanish. If we then fix the gauge A = 0 we recover the original model.



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On the other hand, we can eliminate the non-dynamical gauge fields A, obtaining the dual sigma model.

The existence of global symmetries is a very stringent requirement. A generic metric will not have any Killing vectors.

#### Question

Is it possible to follow the same procedure when the vector fields are not Killing vectors?

# Gauging without isometry

Kotov and Strobl<sup>1</sup> generalised the standard gauging using Lie algebroids (see Kyle's talk). Chatzistavrakidis, Deser, and Jonke<sup>2</sup> applied this non-isometric gauging to

the Buscher procedure we just reviewed.

They introduce a matrix-valued one-form  $\omega_a^b$  satisfying

$$\mathcal{L}_{v_a}g = \omega^b_a \lor \iota_{v_b}g$$
  
 $\mathcal{L}_{v_a}B = \omega^b_a \land \iota_{v_b}B$ 

<sup>1</sup>[1403.8119] <sup>2</sup>[1509.01829] and [1604.03739] The gauged action is almost the same:

$$S_{G}^{\omega} = \int_{\Sigma} g_{ij} DX^{i} \wedge \star DX^{j} + \int_{\Sigma} B_{ij} DX^{i} \wedge DX^{j} + \int_{\Sigma} \eta_{a} F_{\omega}^{a}$$

where the curvature is now given by

$$F_{\omega}^{a} = dA^{a} + \frac{1}{2}C_{bc}^{a}(X)A^{b} \wedge A^{c} - \omega_{bi}^{a}A^{b} \wedge DX^{i}$$

The modified gauge transformations are now

$$\delta_{\epsilon} X^{i} = v_{a}^{i} \epsilon^{a}$$
  
$$\delta_{\epsilon} A^{a} = d\epsilon^{a} + C_{bc}^{a} A^{b} \epsilon^{c} + \omega_{bi}^{a} \epsilon^{b} D X^{i}$$
  
$$\delta_{\epsilon} \eta_{a} = -C_{ab}^{c} \epsilon^{b} \eta_{c} + v_{a}^{i} \omega_{bi}^{c} \epsilon^{b} \eta_{c}$$









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In principle, we could use this to construct T-duals of spaces which have no isometries.

This proposal is equivalent to non-abelian T-duality.<sup>3</sup>

That is, if we can find a set of vector fields and  $\omega_a^b$  which give a non-isometric T-dual, then there exists a set of Killing vectors for the model. The T-dual with respect to these Killing vectors is the same as the non-isometric T-dual.

<sup>3</sup>[1705:09254] P. Bouwknegt, M.B., C. Klimčík, K. Wright

Gauge invariance of the action requires the structure functions to be constant, as well as the vanishing of the following variation:

$$\delta_{\epsilon}(\eta_{a}F_{\omega}^{a}) = \eta_{a}(d\omega_{b}^{a} + \omega_{c}^{a} \wedge \omega_{b}^{c})\epsilon^{b} + \mathcal{O}(A) + \mathcal{O}(A^{2}).$$

We therefore require that  $\omega_a^b$  is flat:

$$R_a^b = d\omega_a^b + \omega_c^b \wedge \omega_a^c = 0,$$

and this tells us that  $\omega_a^b$  is of the form  $K^{-1}dK$  for some  $K_a^b(X)$ .

Using this K, we can perform the following field redefinitions:

$$\begin{aligned} \widehat{A}^{a} &= K_{b}^{a} A^{b} \\ \widehat{\eta}^{a} &= \eta_{b} (K^{-1})_{a}^{b} \\ \widehat{v}_{a} &= v_{b}^{i} (K^{-1})_{a}^{b} \end{aligned}$$

The gauged action can now be rewritten in terms of the new fields  $(X^i, \widehat{A}^a, \widehat{\eta}_a)$ .

$$S_{G}^{\omega}[X,\widehat{A},\widehat{\eta}] = \int_{\Sigma} g_{ij} \widehat{DX}^{i} \wedge \widehat{DX}^{j} + \int_{\Sigma} B_{ij} \widehat{DX}^{i} \wedge \widehat{DX}^{j} + \int_{\Sigma} \widehat{\eta}_{a} \widehat{F}^{a}$$
$$= S_{G}[X,\widehat{A},\widehat{\eta}]$$

where

$$\widehat{F}^{a} = d\widehat{A}^{a} + \frac{1}{2}\widehat{C}^{a}_{bc}\widehat{A}^{b}\wedge\widehat{A}^{c}$$

The gauge transformations become the usual non-abelian gauge transformations, and a short computation reveals

$$\mathcal{L}_{\widehat{v}_a}g = 0$$
  $\mathcal{L}_{\widehat{v}_a}B = 0$ 

#### Conclusion

This proposal is equivalent, via a field redefinition, to the standard non-abelian T-duality

### Examples!

## First example

Consider the 3D Heisenberg Nilmanifold, or twisted torus. It has a metric given by

$$ds^2 = dx^2 + (dy - x \, dz)^2 + dz^2$$



The non-abelian T-dual of this space is given by

$$\widehat{ds^2} = dY^2 + \frac{1}{1+Y^2} \left( dX^2 + dZ^2 \right)$$
$$\widehat{B} = \frac{Y}{1+Y^2} dX \wedge dZ$$

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We can gain a better understanding of the geometry by writing the manifold as a group:

$$Heis := \left\{ \begin{pmatrix} 1 & x & y \\ 0 & 1 & z \\ 0 & 0 & 1 \end{pmatrix} : x, y, z \in \mathbb{R} \right\}$$

(left-invariant) MC forms = (dx, dy - xdz, dz)(right-invariant) vector fields =  $(\partial_x + z\partial_y, \partial_y, \partial_z)$ 

## First example

We could instead try to gauge this space non-isometrically using the left-invariant vector fields:  $\{\partial_x, \partial_y, x\partial_y + \partial_z\}$ .

These are not all isometries:

$$\begin{aligned} \mathcal{L}_{v_1}g &= -dy \otimes dx - dz \otimes dy + 2xdz \otimes dz \\ \mathcal{L}_{v_2}g &= 0 \\ \mathcal{L}_{v_3}g &= dx \otimes dy + dy \otimes dx - xdx \otimes dz - xdz \otimes dx \end{aligned}$$

and they don't commute:

$$[v_1,v_3]=v_2,$$

however...

# First example<sup>4</sup>

If we take  $\omega_3^2 = dx$  and  $\omega_1^2 = -dz$ , with other components vanishing, the non-isometric gauging constraints are satisfied and we can calculate the non-isometric T-dual model.

$$\widehat{ds^2} = dY^2 + rac{1}{1+Y^2} \left( dX^2 + dZ^2 
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<sup>4</sup>Gauged non-isometrically in [1509:01829]

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Unsurprisingly, it is also the T-fold.

A. Chatzistavrakidis, A. Deser, L. Jonke

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# Second example

Consider  $S^3$  with the round metric and B = 0.



This metric has an SO(4) group of isometries, and we can find the non-abelian T-dual with respect to an SU(2) subgroup of this.

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We can write the round metric as

$$g = \lambda^1 \otimes \lambda^1 + \lambda^2 \otimes \lambda^2 + \lambda^3 \otimes \lambda^3$$

where the  $\lambda^i$  are the left-invariant Maurer-Cartan forms.

The right-invariant vector fields are isometries of this metric, so let's try gauging with respect to the left-invariant vector fields<sup>5</sup>.

 $<sup>^5 \</sup>mathrm{These}$  also happen to be isometries of the metric, but let's try to gauge them non-isometrically

The Lie derivatives of the metric with respect to the left-invariant vector fields,  $L_a$  are

$$\mathcal{L}_{L_a}g = -\sum_b C^b_{ac}\lambda^c \lor \lambda^b$$
  
 $= -C^b_{ac}\lambda^c \lor \iota_{L_b}g$ 

We can do non-isometric T-duality by taking  $\omega_a^b = -C_{ac}^b \lambda^c$ .

The remaining gauging constraints are satisfied, and we can calculate the non-isometric T-dual. It is the 'cigar' metric, as expected.



- The equivalence of non-isometric and non-abelian T-duality remains valid for non-exact *H*
- $\bullet$  Geometric interpretation of  $\omega^b_a$  as a connection on a Lie algebroid
- There are proposals for alternate gauging. Unknown how to incorporate into T-duality
  - non-flat  $\omega_a^b$
  - include a term  $\phi^b_{ai} \epsilon^b \star DX^i$  into  $\delta_{\epsilon} A^a$

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#### Thanks!